

IONIC CONDUCTIVITY

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1. THEORY

1.1. Charge transport due to self-diffusion. The *charge transport due to self-diffusion* of α -type molecule λ_{uncorr}^α can be calculated from Einstein relation as

$$(1) \quad \lambda_{uncorr}^\alpha V k_B T = \lim_{t \rightarrow \infty} \frac{\langle q_\alpha^2 [\vec{r}_\alpha(t) - \vec{r}_\alpha(0)]^2 \rangle}{6t},$$

where V and T are volume and temperature of the system, k_B is the Boltzmann constant, $\vec{r}_\alpha(t)$ is vector position at time t for one α -type molecule, and q_α is ionic charge of α -type molecule.

One can improve the statistical precision of results by averaging over the number of α -type molecules in the system N_α ,

$$(2) \quad \lambda_{uncorr}^\alpha = \frac{1}{V k_B T} \lim_{t \rightarrow \infty} \frac{1}{6t} \left\langle \frac{1}{N_\alpha} \sum_{i=1}^{N_\alpha} q_\alpha^2 [\vec{r}_i(t) - \vec{r}_i(0)]^2 \right\rangle.$$

The charge transport due to self-diffusion λ_{uncorr} is defined as

$$(3) \quad \lambda_{uncorr} = \frac{1}{N} \sum_{\alpha=1}^{N_{type}} N_\alpha \lambda_{uncorr}^\alpha,$$

where N_{type} is the number of molecules types, and N is the total number of molecules. Notice that $N = \sum_{\alpha=1}^{N_{type}} N_\alpha$.

1.2. Total charge transport. The *collective charge transport* of α -type and β -type molecules $\lambda^{\alpha\beta}$ can be calculated from Einstein relation as

$$(4) \quad \lambda^{\alpha\beta} V k_B T = \lim_{t \rightarrow \infty} \frac{\langle q_\alpha [\vec{r}_\alpha(t) - \vec{r}_\alpha(0)] \cdot q_\beta [\vec{r}_\beta(t) - \vec{r}_\beta(0)] \rangle}{6t}.$$

One can improve the statistical precision of results by averaging over the number of α -type and β -type molecules in the system N_α and N_β ,

$$(5) \quad \lambda^{\alpha\beta} = \frac{1}{V k_B T} \lim_{t \rightarrow \infty} \frac{1}{6t} \left\langle \frac{1}{N_\alpha} \sum_{i=1}^{N_\alpha} q_\alpha [\vec{r}_i(t) - \vec{r}_i(0)] \cdot \frac{1}{N_\beta} \sum_{j=1}^{N_\beta} q_\beta [\vec{r}_j(t) - \vec{r}_j(0)] \right\rangle.$$

The *total charge transport* λ is defined as

$$(6) \quad \lambda = \frac{1}{N} \sum_{\alpha=1}^{N_{type}} \sum_{\beta=1}^{N_{type}} N_\alpha N_\beta \lambda^{\alpha\beta}.$$

2. IMPLEMENTATION

2.1. Charge transport due to self-diffusion. In realizing Eq. 2 from simulation, the brackets would be interpreted as averages over *time origins*

$$(7) \quad \lambda_{uncorr}^\alpha = \frac{1}{V k_B T} \lim_{t \rightarrow \infty} \frac{1}{6t} \left\{ \lim_{\tau \rightarrow \infty} \frac{q_\alpha^2}{N_\alpha \tau} \sum_{i=1}^{N_\alpha} \int_0^\tau [\vec{r}_i(t_0 + t) - \vec{r}_i(t_0)]^2 dt_0 \right\}.$$

In simulations, one uses discrete time steps, so that one has a set of discrete times $\{t_1, t_2, \dots, t_{N_{tot}}\}$, where $t_i = t_0 + (i - 1)\Delta t$ for $i = 1, 2, \dots, N_{tot}$ and Δt is the *time step*. *Total time simulation* is given by $t_{simul} = t_{N_{tot}} - t_1$. Time origins belong to set of discrete times $\{t_1, t_2, \dots, t_{N_{or}}\}$, where $N_{or} = \frac{N_{tot}}{2}$. *Elapsed time* is set equal to $t = n\Delta t$, where $n = 1, 2, \dots, N_{or}$. Thus,

$$(8) \quad \lambda_{uncorr}^\alpha = \lim_{t \rightarrow \infty} \lambda_{uncorr}^\alpha(t),$$

where

$$(9) \quad \lambda_{uncorr}^\alpha(t) = \frac{e^2}{6t V k_B T} \frac{z_\alpha^2}{N_\alpha N_{or}} \sum_{i=1}^{N_\alpha} \sum_{j=1}^{N_{or}} [\vec{r}_i(t_j + t) - \vec{r}_i(t_j)]^2.$$

Note that $q_\alpha = z_\alpha e$ where e is the elementary charge.

Using Eqs. 8 and 9 into Eq. 3, one can define

$$(10) \quad \lambda_{uncorr} = \lim_{t \rightarrow \infty} \lambda_{uncorr}(t)$$

where

$$(11) \quad \lambda_{uncorr}(t) = \frac{e^2}{6t V k_B T N N_{or}} \sum_{\alpha=1}^{N_{type}} z_\alpha^2 \sum_{i=1}^{N_\alpha} \sum_{j=1}^{N_{or}} [\vec{r}_i(t_j + t) - \vec{r}_i(t_j)]^2.$$

From *apparent* self-diffusion constant of α -type molecule $D_\alpha(t)$, Eq. 11 can be written as

$$(12) \quad \lambda_{uncorr}(t) = \frac{e^2}{Vk_B T} \sum_{\alpha=1}^{N_{type}} z_\alpha^2 n_\alpha D_\alpha(t),$$

where $n_\alpha = N_\alpha/N$.

2.2. Total charge transport. In realizing Eq. 5 from simulation, the brackets would be interpreted as averages over *time origins*

$$(13) \quad \lambda^{\alpha\beta} = \frac{1}{Vk_B T} \lim_{t \rightarrow \infty} \frac{1}{6t} \left\{ \lim_{\tau \rightarrow \infty} \frac{q_\alpha q_\beta}{N_\alpha N_\beta \tau} \sum_{i=1}^{N_\alpha} \sum_{j=1}^{N_\beta} \int_0^\tau [\vec{r}_i(t_0 + t) - \vec{r}_i(t_0)] \cdot [\vec{r}_j(t_0 + t) - \vec{r}_j(t_0)] dt_0 \right\}.$$

In simulations, one uses discrete time steps, so that one has a set of discrete times $\{t_1, t_2, \dots, t_{N_{tot}}\}$, where $t_i = t_0 + (i - 1)\Delta t$ for $i = 1, 2, \dots, N_{tot}$ and Δt is the *time step*. *Total time simulation* is given by $t_{simul} = t_{N_{tot}} - t_1$. Time origins belong to set of discrete times $\{t_1, t_2, \dots, t_{N_{or}}\}$, where $N_{or} = \frac{N_{tot}}{2}$. *Elapsed time* is set equal to $t = n\Delta t$, where $n = 1, 2, \dots, N_{or}$. Thus,

$$(14) \quad \lambda^{\alpha\beta} = \lim_{t \rightarrow \infty} \lambda^{\alpha\beta}(t),$$

where

$$(15) \quad \lambda^{\alpha\beta}(t) = \frac{e^2}{6tVk_B T} \frac{z_\alpha z_\beta}{N_\alpha N_\beta N_{or}} \sum_{i=1}^{N_\alpha} \sum_{j=1}^{N_\beta} \sum_{k=1}^{N_{or}} [\vec{r}_i(t_k + t) - \vec{r}_i(t_k)] \cdot [\vec{r}_j(t_k + t) - \vec{r}_j(t_k)].$$

Using Eqs. 14 and 15 into Eq. 6, one can define

$$(16) \quad \lambda = \lim_{t \rightarrow \infty} \lambda(t)$$

where

$$(17) \quad \lambda(t) = \frac{e^2}{6tVk_B T N N_{or}} \sum_{\alpha=1}^{N_{type}} z_\alpha \sum_{\beta=1}^{N_{type}} z_\beta \sum_{i=1}^{N_\alpha} \sum_{j=1}^{N_\alpha} \sum_{k=1}^{N_{or}} [\vec{r}_i(t_k + t) - \vec{r}_i(t_k)] \cdot [\vec{r}_j(t_k + t) - \vec{r}_j(t_k)].$$

For practical reasons, one defines the *degree of uncorrelated ion motion* as

$$(18) \quad \gamma_d = \lim_{t \rightarrow \infty} \gamma_d(t),$$

where

$$(19) \quad \gamma_d(t) = \frac{\lambda(t)}{\lambda_{uncorr}(t)}.$$

Then, one calculates $\lambda_{uncorr}(t)$ from Eq. 12 using diffusion coefficients corrected for the finite simulation cell size as well as *apparent* degree of uncorrelated ion

motion $\gamma_d(t)$. Finally, *apparent* total charge transport $\lambda(t)$ is obtained from Eq. 19.